Progress in self-consistent evolution with a time domain scalar charge self-force code

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June 26, 2018 21st Capra Meeting 2018, AEI, Potsdam, Germany We wish to determine the self-forced motion and field (e.g. energy and angular momentum fluxes) of a particle with scalar charge

$$\Box \psi^{\text{ret}} = -4\pi q \int \delta^{(4)}(x - z(\tau)) \, d\tau.$$

Two general approaches:

- Compute enough "geodesic"-based self-forces and then use these to drive the motion of the particle. (Post-processing, fast, accurate self-forces, relies on slow orbit evolution)
- Compute the "true" self-force <u>while</u> simultaneously driving the motion. (Potentially slow and expensive, potentially less accurate self-forces)

## Effective source approach.

... is a general approach to self-force and self-consistent orbital evolution that doesn't use any delta functions.

Key ideas

 $\blacktriangleright$  Compute a regular field,  $\psi^{\rm R},$  such that the self-force is

$$F_{\alpha} = \nabla_{\alpha} \psi^{\mathsf{R}}|_{x=z},$$

where  $\psi^{\mathsf{R}} = \psi^{\mathsf{ret}} - \psi^{\mathsf{S}}$ , and the Detweiler-Whiting singular field  $\psi^{\mathsf{S}}$  can be approximated via local expansions:  $\psi^{\mathsf{S}} = \tilde{\psi}^{\mathsf{S}}(x|z, u, a) + O(\epsilon^n)$ .

▶ The effective source, S, for the field equation for  $\psi^{\mathsf{R}}$  is regular at the particle location

$$\Box \psi^{\mathsf{R}} = \Box \psi^{\mathsf{ret}} - \Box \tilde{\psi}^{\mathsf{S}} = S(x|z, u, a, \dot{a}, \ddot{a}),$$

where  $\Box \tilde{\psi}^{\mathsf{S}} = -4\pi q \int \delta^{(4)}(x-z(\tau)) \, d\tau - S.$ 

# Self-consistent vs. geodesic evolutions.

- One main goal is to compare our self-consistent evolutions with Niels Warburton's geodesic evolutions.
- First attempt: 3+1D multi-patch finite difference code with a  $C^0$  effective source.
- ► 3+1D accuracy limited by the non-smoothness of the source leading to high frequency noise with 2nd order convergent amplitude.
- Self-consistent evolutions agreed beautifully with geodesic evolutions within the errors (dominated by the noise).
- Next attempt: Improve the effective source smoothness to  $C^2$ .
- ► Geodesic evolution agreed with the C<sup>0</sup> evolutions and the frequency domain result with the noise reduced by more than an order of magnitude.
- ► However, we found differences between C<sup>2</sup> and C<sup>0</sup> results as soon as the back-reaction was turned on.
- 1+1D discontinuous Galerkin code without acceleration terms lost mode sum convergence when back-reaction was turned on.
- Now: 1+1D discontinuous Galerkin code with acceleration terms in the effective source.

# The code.

- Solves the spherical harmonic decomposed scalar wave equation in a Schwarzschild spacetime with a scalar effective source.
- ▶ Uses the Discontinuous Galerkin method for spatial discretization.
- ▶ Uses the method of lines approach and supports a number of time integrators.
- ► Uses a world-tube approach.
- ► Uses hyperboloidal slices, placing the computational domain boundaries at the horizon and 𝒴<sup>+</sup>.
- Uses a time dependent coordinate transformation to place the particle at a fixed coordinate location.
- The effective source include acceleration terms.
- $\blacktriangleright$  Can read in frequency domain code initial data for small  $\ell$  modes.



## The state of self-consistent evolution at last Capra.

p = 9.9, e = 0.1, q = 1/8. Only four-acceleration passed in to the effective source!!!!



## What was wrong?

- Noticed that a bit of noise appeared in the extracted self-force shortly after each periapsis passage.
- ▶ Noticed that the same thing happened after each apapsis passage.
- Turns out it was caused by the calculation of  $u^r$ .

$$u^{r} = \pm \sqrt{E^{2} - \left(1 - \frac{2M}{r}\right)\left(1 + \frac{L_{z}^{2}}{r^{2}}\right)},$$

$$E^2 = \frac{(p-2-2e)(p-2+2e)}{p(p-3-e^2)} \text{ and } L^2_z = \frac{p^2 M^2}{p-3-e^2}$$

- Instead of  $u^r \approx 10^{-16}$  we got  $u^r \approx 10^{-8}$  just before and after peri- and apapsis.

- This apparently generates enough noise to trigger a feedback instability when the back-reaction is applied.
- Easy fix: do this calculation in quad precision.

p = 9.9, e = 0.1, q = 1/8. Still only four-acceleration passed in to the effective source.



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- We still have instabilities when we pass in  $\dot{a}$  and/or  $\ddot{a}$ .
- At first we thought this was caused by extra noise in the numerical calculation of time derivatives of the 4-acceleration.
- ► We experimented with smoothing finite differences (made complicated due to time varying ∆t): Did not help much.
- We then implemented an Adams-Bashford-Moulton multi-value (ABMV) time integrator where higher time derivatives of the variables are part of the evolution system: Extended the run time but still instabilities.
- ▶ Question is: How important are the higher derivatives of *a*?
- It turns out that the ABMV scheme does allow for long evolutions if lmax is not too large. Comparing the phase evolution between runs with and without à terms may help quantify this.



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## Conclusions and Outlook.

- We can now do self-consistent orbits where the effective source depends on the acceleration.
- ▶ Preliminary results consistent with the expectation that the phase error between 'geodesic' and 'self-consistent' evolutions grows as  $(t t_0)^2$  and scales as  $q^2$ .
- Need to understand the numerical errors better before we can say anything definite about the magnitude of the phase error.
- Need to finish investigation into importance of time derivatives of the acceleration. REU student Aaron Hodson is working on this.
- Gravitational perturbation codes (Lorenz, Regge-Wheeler-Zerilli and Teukolsky) are in various stages of development/testing.
- Currently undertaking a redesign and rewrite of the code. Once this is done, we plan to release the code as open software (http://bhptoolkit.org?).
- Plan to extend the code to handle Kerr as well.

#### Teaser: New code for Teukolsky in Schwarzschild by Sarah Skinner

